

Densities and temperatures in the Venus mesosphere and lower thermosphere retrieved from SOIR on board Venus Express.

Carbon dioxide measurements at the Venus terminator.

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Abstract

SOIR is a high resolution spectrometer flying on board the ESA Venus Express mission. It performs solar occultations of the Venus high atmosphere, and so defines unique vertical profiles of many of the Venus key species. In this paper, we focus on the Venus main constituent, carbon dioxide. We explain how the temperature, the total density and the total pressure are derived from the observed CO₂ density vertical profiles. A striking permanent temperature minimum at 125 km is observed. The dataset is processed in order to obtain a Venus Atmosphere from SOIR measurements at the Terminator (VAST) compilation for different latitude regions and extending from 70 up to 170 km in altitude. The results are compared to many literature results obtained from ground-based observations, previous missions, and the Venus Express mission. The homopause altitude is also determined.

1 Introduction

The Venus terminator in the mesosphere and thermosphere is a region of great interest as a transition between the hot dayside and the cold nightside of the Venus atmosphere [Keating *et al.*, 1980]. It was previously supposed that a compression of the atmosphere at these locations might induce large temperature variations, due to the same pressure levels altitude difference between the day and the night sides of the planet [Parkinson, 2012]. However, these are poorly

29 known, since only a few measurements were obtained at these particular local solar times, i.e. 6.00 AM and 6.00 PM.
30 The SOIR instrument on board Venus Express (VEX) offers the possibility to routinely monitor the density of various
31 key Venus species and temperature vertical profiles [Bertaux *et al.*, 2007b; Fedorova *et al.*, 2008; Vandaele *et al.*, 2008].
32 SOIR is particularly well suited to help identify the dynamic and chemical processes that take place in this region.

33

34 Carbon dioxide is the main component of Venus' atmosphere, with a mean volume mixing ratio (VMR) of 96.5 % up to
35 an altitude of 120 km [Keating *et al.*, 1980]. Above this, its VMR decreases with altitude as CO₂ is photodissociated on
36 the dayside of the planet by solar ultraviolet radiation, forming carbon monoxide and excited oxygen atoms. As the main
37 constituent of the Venus atmosphere, carbon dioxide is of great interest in order to describe the physics of the
38 atmosphere at the terminator. The SOIR instrument measures CO₂ absorption across a broad spectral window. The
39 observed atmospheric transmittance spectra are subsequently inverted to obtain vertical density profiles, which extend
40 from 70 km up to 170 km. This paper presents the CO₂ density and temperature profiles of selected orbits that were
41 obtained between 2006 and 2011, and compiles them in order to present a global view of the atmospheric characteristics
42 at the Venusian terminator.

43

44 Previously, CO₂ density and temperature have been measured in this particular altitude range and at, or close to, the
45 terminator in the Venus atmosphere from the ground [Clancy *et al.*, 2003; Clancy *et al.*, 2008; 2012; Lellouch *et al.*,
46 1994] (sub-mm range from Earth telescopes) as well as from previous spacecrafts and in-situ probes: Pioneer Venus
47 [Fox and Kasprzak, 2007; Keating *et al.*, 1980; von Zahn *et al.*, 1980 and references therein] and VEX instruments
48 [Bertaux *et al.*, 2007b; Irwin *et al.*, 2008; Piccialli *et al.*, 2012; Tellmann *et al.*, 2009]. Models have been developed
49 based on previous measurement data sets obtained during the Pioneer Venus era, such as the VTS3 model of the upper
50 atmosphere [Hedin *et al.*, 1983], the Keating model of the upper atmosphere [Keating *et al.*, 1980] and the Zasova's
51 model from the surface up to 100 km [Zasova *et al.*, 2007; Zasova *et al.*, 2006]. Altogether, these measurements and
52 models cover an altitude range between 50 and 200 km, and contain information encompassing the terminator region.
53 Large temporal variations are observed in these data sets.

54 **2 Instrument description and observation geometry**

55 The SOIR instrument has been extensively described in Bertaux *et al.* (2007a), Nevejans *et al.* (2006), Mahieux *et al.*,
56 (2008), Mahieux *et al.* (2009) and Mahieux *et al.* (2010). We will summarize here only the most important SOIR
57 characteristics.

58

59 SOIR is one of the three channels of the SPICAV/SOIR instrument. It is an infrared spectrometer, using an echelle
60 grating as the diffracting element. The accessible wavenumber range covers the 2200 to 4400 cm⁻¹ region, and is divided

61 into 94 useful diffracting orders, or simply orders, from 101 to 194. The order selection is performed using an acousto-
62 optic tunable filter (AOTF), which allows to select and transmit only a small wavenumber range diffracted by the echelle
63 grating. The resolution of SOIR varies from order to order, with a value of about 0.11 cm^{-1} in order 101 to 0.21 cm^{-1} in
64 order 194. The SOIR detector is composed of 320 pixels in its spectral direction. The spectral width of a pixel varies
65 from 0.06 to 0.12 cm^{-1} and the free spectral range (FSR) has a constant value of 22.4 cm^{-1} . The SOIR useful detector
66 pixels are combined into two groups in the spatial direction called bins. Two simultaneous measurements are thus
67 obtained at two slightly different altitudes corresponding to the 2 bins on the detector. The width of the AOTF bandpass
68 transfer function is $\sim 24 \text{ cm}^{-1}$, while the width of an order varies between 19.3 cm^{-1} to 37.1 cm^{-1} with ascending order,
69 which causes an order overlapping on the detector. Moreover, the AOTF transfer function is not zero in the adjacent
70 orders [Mahieux *et al.*, 2009]. To correctly simulate the SOIR measurement, more than one diffraction order have to be
71 taken into account [Mahieux *et al.*, 2010]. Five contiguous diffraction orders are usually simulated, i.e. the scanned one
72 and two adjacent orders on each side.

73

74 SOIR performs solar occultation observations of the Venus atmosphere from the VEX spacecraft, which is in a polar
75 orbit with its periapsis located above the North Pole. The vertical size of the instantaneously scanned atmosphere at the
76 limb tangent point varies from a few hundreds of meters for the Northern measurements to tens of kilometers for the
77 Southern measurements. The altitude range probed by SOIR, i.e. where measurements are scientifically meaningful,
78 varies from 70 km up to 170 km. The lower boundary corresponds to total absorption of sunlight by Venus' clouds, and
79 the upper boundary to the detection of the strongest CO_2 band in the selected SOIR wavenumber range. During an
80 occultation, four different diffraction orders are measured quasi-simultaneously – four sequentially within 1 s, each one
81 lasting 160 ms at maximum. It allows us to study either the same species at different ranges of altitudes, as it will be the
82 case here for CO_2 , or different species to obtain volume mixing ratios [Bertaux *et al.*, 2007b; Fedorova *et al.*, 2008;
83 Vandaele *et al.*, 2008]. The occultations are grouped in occultation seasons (OS), which are time periods when the solar
84 occultations take place relative to the VEX spacecraft orbital configuration. These OS periods occur roughly every three
85 months for one month. The definition of the OS is a parameter linked to the VEX spacecraft orbital characteristics, and
86 has no sense in terms of Venus climatology.

87

88 The spectroscopic parameters are obtained from the Hitran 2008 database [Rothman *et al.*, 2009], with corrected values
89 for the pressure broadening coefficients and shifts, to take into account that the atmosphere of Venus is mainly CO_2
90 instead of nitrogen and oxygen [Vandaele *et al.*, 2008]. Carbon dioxide is studied in the range 2400 to 4000 cm^{-1} (2.5 to
91 $4.16 \mu\text{m}$). In this study we essentially concentrate on the vibrational bands of the following five isotopologues $^{12}\text{C}^{16}\text{O}_2$,
92 $^{13}\text{C}^{16}\text{O}_2$, $^{16}\text{O}^{12}\text{C}^{18}\text{O}$, $^{16}\text{O}^{13}\text{C}^{18}\text{O}$ and $^{16}\text{O}^{12}\text{C}^{17}\text{O}$. The isotopic ratio of the carbon and oxygen atoms is assumed to be the

93 same as on Earth [Bézard *et al.*, 1987; Clancy and Muhleman, 1991]. For $^{12}\text{C}^{16}\text{O}_2$, the 20011-10002 and 00021-01101
94 bands (in Hitran notation [Rothman *et al.*, 2009]) are observed in orders 171 (3822 to 3854 cm^{-1}) and 176 to 180 (3933
95 to 4057 cm^{-1}), respectively. They give information on low altitudes (70 to 90 km). The 21103-00001 band is measured in
96 order 141 (3151 to 3178 cm^{-1}) and the 21102-00001 band in orders 147 to 150 (3285 to 3381 cm^{-1}); they correspond to
97 mid altitudes (80 to 120 km). Finally the 10012-00001 and 10011-00001 bands are measured in the orders 160 to 168
98 (3576 to 3787 cm^{-1}) and are observed at very high altitudes from 120 to 170 km. The 10012-00001 band of $^{13}\text{C}^{16}\text{O}_2$ is
99 measured in orders 155 and 156 (3464 to 3516 cm^{-1}); it is measured at high altitude from 100 to 140 km. Three bands of
100 $^{16}\text{O}^{12}\text{C}^{18}\text{O}$ are observed, the 01111-00001 band in order 134 (2995 to 3020 cm^{-1}) at low altitude (from 70 to 90 km), and
101 the 20002-00001 and 20003-00001 bands at mid-altitude (90 to 120 km) measured in orders 112 (2503 to 2525 cm^{-1}),
102 and 117 and 118 (2615 to 2637 cm^{-1}) respectively. The 20001-00001 band of the $^{16}\text{O}^{13}\text{C}^{18}\text{O}$ isotopologue is observed in
103 order 121 (2704 to 2727 cm^{-1}) at low altitude, from 70 to 90 km.

104 **3 Retrieval method description**

105 The core of the algorithm used in this paper, Asimat, has been described previously in the companion paper [Mahieux *et al.*
106 *et al.*, 2010]. However, as it is presented and explained hereafter, the method is completely revised in terms of temperature
107 retrieval. Only a short summary of what is taken from the previously published method will be given here. The density
108 and the background parameters are retrieved simultaneously from the solar occultation data, using the Optimal
109 Estimation (OE) method developed by Rodgers [Rodgers, 1990; 2000]. The background parameters are fifth-order
110 polynomial expressions that fit the backgrounds of each spectrum. They account for the absorption of all non-molecular
111 species, like aerosols [Wilquet *et al.*, 2012], or any molecule that has a structure broader than a SOIR diffraction order.

112

113 The useful altitude range varies depending on the order studied and the targeted species. The maximum altitude
114 corresponds to the tangent altitude at which absorption structures appear in the spectra. The minimum altitude is
115 obtained when the absorption lines are saturated. The same approximation in terms of inversion scheme as in the onion
116 peeling method is used to discretize the atmosphere: the atmosphere is assumed to be spherically homogeneous at each
117 altitude level [Mahieux *et al.*, 2010; Vandaele *et al.*, 2008]. This is an approximation, as it is known that the atmosphere
118 undergoes density variations across the terminator, with larger densities on the day side than the night side at same
119 altitude levels. However, the steepness and the shape of these gradients are currently unknown. Attempts have been
120 made to introduce a horizontal dependency of the density orthogonally to the terminator. This does not change the shape
121 of the retrieved vertical density profiles, but only scales the whole density profile by a factor proportional to the
122 introduced dependency. For this reason, the density gradient is not taken into account in the following, though it will be
123 further investigated in the future.

124

125 In the companion paper *Mahieux et al. (2010)*, we claimed that the rotational temperature could be derived directly from
126 the spectra, using the rotational structure of the CO₂ absorption bands which are resolved. However, it has been observed
127 that this approach leads to incorrect temperature profiles [*Mahieux et al., 2011*]. The envelope of the rotational structure
128 as seen in the measured spectra is modulated by the instrument AOTF bandpass function during the acquisition process.
129 The shape of this function is unfortunately slightly broader to the rotational structure envelope, and obscures the
130 dependency of the measured rotational envelope on gas rotational temperature. The retrieved temperature profiles
131 obtained using the Rodgers method were unchanged from the a-priori temperature profiles, implying that no information
132 on the local temperature could be gained from measurements. In addition, density and temperature profiles resulting
133 from the inversion did not verify hydrostatic equilibrium, which is not physically valid. Altogether, this clearly indicates
134 that the temperature profiles are not directly retrievable, using the rotational information contained in the spectra.

135

136 In order to overcome this issue, a global iterative approach is now considered, in which each iteration is performed in
137 three steps, see the diagram in Figure 1. First, the CO₂ density is calculated using the previous iteration temperature
138 profile, then a new temperature profile is obtained from the new CO₂ density profile: at each iteration, the Asimat
139 algorithm is applied with the temperature profile fixed, on all the simultaneous sets of spectra from all orders and bins in
140 which CO₂ is measured. In a second step, the CO₂ density profiles are combined in a single one, and in a third step the
141 hydrostatic equilibrium equation is used to derive the new temperature profile.

142

143 In order to illustrate the method, let us consider an orbit for which the four diffraction orders are dedicated to CO₂
144 measurements which cover the whole altitude range, i.e. between 85 and 155 km of altitude. The Asimat algorithm is run
145 on the 8 sets of spectra (one spectrum per bin per order). The initial temperature profile is chosen as the Zasova model
146 [*Zasova et al., 2007; Zasova et al., 2006*] below 100 km and the Keating model [*Keating et al., 1980*] above 140 km.
147 The inversions result in a set of eight CO₂ density vertical profiles. These independent density profiles are combined into
148 a single one, using a moving linear mean square regression, weighted by the error of each individual density value and
149 that considers all the points located within one scale height of altitude. One single smooth CO₂ recombined density
150 profile is thus obtained. SOIR typical vertical sampling varies with the latitude of the measurement: 1 km above 70°N,
151 500 m between 40° and 70°N, and from 1 to 4 km south of 40°N, with the largest values at the South Pole. Tangent
152 altitudes vary slightly among the eight data set spectra requiring interpolation of the atmospheric results on a new
153 altitude scale. An altitude step of 1 km has been chosen. Smaller altitude steps were also considered, but this step did not
154 influence the results.

155

156 Figure 1

157

158 At the end of each iteration, a new temperature profile is calculated from the hydrostatic equation and the ideal gas law
159 using the CO₂ vertical profile:

$$\begin{cases} dp(z) = -\rho(z) \cdot MM(z) \cdot g(z) \cdot dz \\ g(z) = g_0 \cdot \left(\frac{R_{Venus}}{R_{Venus} + z} \right)^2 \\ p(z) = \rho(z) \cdot k_B \cdot T(z) \\ \rho_{CO_2}(z) = VMR_{CO_2}(z) \cdot \rho(z) \end{cases} \quad (1)$$
$$\Rightarrow T(z) = -\frac{VMR_{CO_2}(z) \cdot MM_{CO_2}}{k_B \cdot \rho_{CO_2}(z)} \cdot \int_{z_0}^z \frac{(\rho_{CO_2}(\tilde{z}) - \rho_{CO_2}(z_0)) \cdot g(\tilde{z})}{VMR_{CO_2}(z)} \cdot d\tilde{z}$$

160 where z is the altitude, z_0 is the reference altitude, $p(z)$ is the total pressure, $\rho_{CO_2}(z)$ is the CO₂ volume density and
161 $\rho(z)$ is the total volume density, $VMR_{CO_2}(z)$ is the CO₂ VMR, MM_{CO_2} is the CO₂ molar mass (44.01 g/mol) and
162 $MM(z)$ is the atmosphere molar mass, $g(z)$ is the Venus gravitational acceleration, g_0 is the Venus gravitational
163 constant (8.87 m/s²), R_{Venus} is the Venus radius (6053.1 km), $T(z)$ is the temperature profile and k_B is the Boltzmann
164 constant (1.38 · 10⁻²³ J/K). The altitude step Δz chosen in the altitude scale equals to 1 km. The reference altitude z_0 is
165 always taken as the altitude of the highest retrieved point. The CO₂ VMR is taken from the Keating [Keating *et al.*,
166 1980] and Zasova [Zasova *et al.*, 2007; Zasova *et al.*, 2006] models, as a function of altitude.

167

168 Because CO₂ absorption cross sections (ACS) are temperature dependent, the Asimat algorithm is iteratively run on the
169 8 sets of measured spectra, considering the last temperature profile obtained. The convergence is assumed to be achieved
170 when both recombined density and temperature profiles stabilize. In the present study we specify that the temperature
171 difference between two successive iterations should be less than 7 K, and the density natural logarithm difference less
172 than 0.3, both calculated as integrated difference on the whole profile. The evolution of those two convergence criteria is
173 shown in Figure 2 for a typical observation (in this case, orbit 687.1 – 03/08/2008). The total root mean square (RMS)
174 value obtained by summing all residuals of all the spectra included in the retrieval (all orders, all bins, and all altitudes)
175 is also shown in Figure 2. Considering the convergence criteria on the density and the temperature and allowing for the
176 total RMS to stabilize to a low value, it can be seen that in the illustrated example, convergence is achieved after 3 steps.
177 Typically, convergence is reached after 4 steps (3.93 ± 1.44). Figure 3 illustrates the evolution of the density and the
178 temperature profiles during the iterative process. The density profile departs from initialization notably at 120-130 km
179 altitudes, associated with a ~100 K variation in temperature over this altitude region.

180

181 Figure 2

182

183 Figure 3

184

185 The hydrostatic equation is applied in the described method whether the measurements are taken above or below the
186 homopause, since Eq. (1) considers the total density of the atmosphere. The homopause is located at an altitude of about
187 125 km on Venus or at a pressure level of $2 \cdot 10^{-5}$ mbar [de Pater and Lissauer, 2001]. Below that level, in the
188 homosphere, each species adopts the same scale height, which is given by the average molecular mass of the
189 atmosphere, whereas above, in the heterosphere, molecular diffusion dominates over vertical mixing [Jacobson, 1999];
190 each individual species follows its own scale height.

191

192 From Eq. (1), it can be seen that the temperature profile is dependent on the density value derived at the top of the
193 profile: $\rho_{CO_2}(z_0)$. In order to study the influence of this parameter, the temperature profile has been calculated for
194 density values at the top the atmosphere varying between $\rho_{CO_2}(z_0) - 2 \cdot \sigma$ and $\rho_{CO_2}(z_0) + 2 \cdot \sigma$, with σ the density
195 error at the top of the profile. The results are presented in Figure 4. The impact of the starting density remains below the
196 retrieved temperature error bars, which indicates that the influence of this parameter does not affect much the
197 temperature retrieval. However, the temperature profile is strictly reliable 10 km below the maximum altitude, which is
198 located between 1 and 2 scale heights above the homopause altitude. The temperature profiles are robust below 140 km,
199 while larger error bars are observed at higher altitude.

200

201 Figure 4

202

203 Finally, knowing the temperature profile derived using the method described just above, and considering the CO₂ VMR
204 from the models of Keating and Zasova [Keating et al., 1985; Zasova et al., 2007; Zasova et al., 2006], which is the only
205 information that could not be directly retrieved from the SOIR measurements, it is possible to derive the total
206 atmospheric pressure and density and the scale height, $H(z)$:

$$\begin{cases} \rho(z) = \frac{\rho_{CO_2}(z)}{VMR_{CO_2}(z)} \\ p(z) = \rho(z) \cdot k_B \cdot T(z) \\ H(z) = \frac{k_B \cdot T(z)}{g(z) \cdot MM_{CO_2} \cdot VMR_{CO_2}(z)} \end{cases} \quad (2)$$

207 The converged CO₂ density and temperature profiles for two typical orbits (687.1 and 1567.1) are given in Figure 5
208 along with their respective typical errors. A North polar and an equatorial measurement are chosen, in order to show the

209 influence of the vertical sampling on the profiles – see Table 1 for the exact location of the measurements. The vertical
210 sampling of the polar orbit is much larger than for the equatorial orbit, around 500 m and 3 km respectively. In the top
211 part of the figure, the two panels present the individual CO₂ density values covering the altitudes from 80 to 160 km,
212 along with the averaged CO₂ vertical profile. The error level on the density for the two presented orbits is between 1 and
213 5 %. The temperature profiles, in the bottom panels, are obtained from the CO₂ density profile using Eq. (1), assuming
214 the hydrostatic equilibrium and the CO₂ VMR from the Keating and Zasova models. The error on the temperature, given
215 by the gray envelopes, varies typically from 1 to 20 K and is higher when the slope in the CO₂ density varies rapidly or
216 at high altitude where the CO₂ VMR has smaller values. The cold layer at the altitude of 110-140 km results from the
217 change of curvature of the CO₂ density profile at the same altitude level.

218

219 Figure 5

220

221 **4 Orbit selection and localization**

222 Although the SOIR dataset now contains 478 observations (up to Jan. 2012) from which 465 can be used to retrieve
223 information on CO₂, only a subset has been considered in this study to define the Venus Atmosphere from SOIR
224 measurements at the Terminator (VAST). Indeed we only considered the orbits for which CO₂ has been measured at
225 least once at very high altitude, i.e. in which the strong ¹²C¹⁶O₂ absorption band at 3 μm is present. This corresponds to
226 SOIR orders 160 to 166. This band is usually observed from an altitude as high as 165 km. Using this criteria, only 59
227 observations were considered. The selected measurements are obtained on AM or PM terminator sides for a wide range
228 of latitudes. In terms of time coverage, they were obtained during various occultation seasons (from OS 1 to OS 17)
229 between 2006 and 2011. Figure 6 shows the localization in terms of orbit number, latitude and local solar time of the
230 measurements. The fact that the local solar times displayed in Figure 6 are not equal to 6.00 AM or 6.00 PM when
231 approaching to the pole is an artifact, and comes from the fact that the notion of local solar time becomes meaningless at
232 these latitudes. This means that the only local solar time information that should be used is the terminator side, either
233 AM or PM. Details of the selected orbits are summarized in Table 1. In the future, VAST developed here will be used as
234 starting condition for the resolution of Eq. (1). The aim is to gradually incorporate new retrievals as they are obtained to
235 refine the model.

236

237 Table 1

238

239 Due to the VEX orbit, two successive measurements, which occur on a 24-Earth hour basis, are taken at somewhat
240 different latitudes on the same side of the terminator: there is a time-latitude gap between successive measurements, such

241 that it is not easy to discriminate between time and latitude variations of the density or temperature profiles. In
242 consequence, latitudinal and short term trends should be considered with care, as they are intrinsically linked.

243

244 Figure 6

245

246 Short term CO₂ density variations for a given latitude region of the terminator are not considered in the current study,
247 since they are removed by the use of a statistically large enough sample of profiles. Systematic long term CO₂ density
248 variations are not observed in the subset studied here. These variations could obviously influence the model, and will be
249 investigated further in the future.

250 **5 Carbon dioxide density and temperature profiles**

251 The carbon dioxide density and temperature profiles are presented in Figure 7 for the whole set. The error bars are not
252 displayed to ensure readability of the figure.

253

254 Figure 7

255

256 The density profiles all show the same shape, with a small yet systematic latitudinal dependency. They present a change
257 of slope – a curvature – in the logarithmic scale around 120-140 km of altitude. The steepness of this gradient may
258 change slightly from one orbit to another, and appears to also depend on latitude. Also, Figure 7 indicates that a global
259 variability of the CO₂ density profiles (within a factor 10) is observed as a function of time and/or latitude. At constant
260 density, the variations are equivalent to 2 scale heights (the scale height H is approximately 3 to 5 km). The general
261 relative error on the density profiles is 1 % to 10 %, except in the region of the curvature change (120 km-140 km)
262 where larger errors are observed, 10 % to 40 % . .

263

264 The temperature profiles are given as a function of total pressure to ensure consistent comparison between the different
265 curves, in other terms to remove the influence of the observed local variations in the CO₂ density. The change of slope in
266 the density profiles observed in Figure 7 corresponds to a temperature minimum in the 10⁻⁵ mbar region (120-130 km of
267 altitude), with low temperatures between 60 and 110 K. This minimum is surrounded by two temperature maxima,
268 located in the 5·10⁻⁷ mbar region or 130-140 km of altitude range (200-350 K) and the 5·10⁻³ mbar region or 100-110 km
269 of altitude range (180-250 K). For a given pressure level, the temperature profiles also show variability of the order of
270 50 K, at all altitudes. At constant temperature, the variations are also equivalent to 2 scale heights (around 3 to 5 km).
271 The error bars for the temperature profiles vary between 2 and 10 K below 120 km and between 10 and 60 K above
272 120 km. These values are directly linked to the error on the density profiles and to the value of the CO₂ VMR obtained

273 from the Keating and Zasova models, through the hydrostatic equilibrium in Eq. (1). Larger temperature errors are
 274 observed at high altitude due to the lower values of the CO₂ VMR.

275

276 The CO₂ density profiles and the temperature profiles are compared in Figure 7 to the Keating and Zasova models. In
 277 terms of density, the agreement is good at altitudes lower than 120 km, but the model does not reproduce the observed
 278 CO₂ density drop measured by the SOIR instrument. This is reflected in the temperature profiles, as the measured cold
 279 layer in the 10⁻⁵ mbar region is not present in the Keating and Zasova models [Keating *et al.*, 1980; Zasova *et al.*, 2007;
 280 Zasova *et al.*, 2006]. At higher pressure level (lower altitudes), the agreement between the model and the SOIR
 281 temperatures is better, but differences can reach up to 50 K.

282 6 CO₂ homopause altitude

283 The homopause is the altitude level at which the molecular diffusion and eddy diffusion processes are of the same order
 284 of magnitude. Above the homopause, each species assumes its own scale height based on its own molecular mass. The
 285 composition of the atmosphere varies with altitude even in the homosphere, the difference with the heterosphere being
 286 that all species share the same hydrostatic scale height, computed from the local mean molecular mass.

287

288 The molecular diffusion coefficient of CO₂ (D_{CO_2}) is obtained using the relations [Jacobson, 1999]:

$$D_{CO_2} = \frac{3 \cdot 10^{-2} \cdot \pi \cdot \sqrt{2}}{16} \cdot l_{CO_2} \cdot v_{CO_2}^{th} \text{ with } l = \frac{1}{Q_{CO_2} \cdot \rho_{CO_2}} \text{ and } v_{CO_2}^{th} = \sqrt{\frac{3 \cdot k_B \cdot T \cdot N_{Av}}{MM_{CO_2}}} \quad (3)$$

289 where l_{CO_2} is the free mean molecular path, $v_{CO_2}^{th}$ the mean molecular thermal speed, Q_{CO_2} the CO₂ effective cross-
 290 section (0.52 nm²) and N_{Av} the Avogadro number (6.02 · 10²³ molec).

291

292 The eddy coefficient is calculated using [Brecht *et al.*, 2012; von Zahn *et al.*, 1979; von Zahn *et al.*, 1980]:

$$K = \frac{5.5 \cdot 10^{12}}{\sqrt{\rho}} \quad (4)$$

293 The homopause level corresponds to the altitude where the molecular diffusion coefficient equals the eddy diffusion
 294 coefficient [Jacobson, 1999]. The derived homopause altitude varies between 124 and 134 km as shown in Figure 8.
 295 This is in agreement with the literature [de Pater and Lissauer, 2001; Jacobson, 1999; von Zahn *et al.*, 1980]. For
 296 example, de Pater and Lissauer, 2001 reported values at 2 · 10⁻⁵ mbar, which corresponds to 120 to 132 km. A large
 297 altitude variability is observed, corresponding to two scale heights, but no latitudinal trend.

298

299 Figure 8

300

301 **7 Venus Atmosphere from SOIR measurements at the Terminator (VAST)**

302 In order to study the global distribution of the latitudinal parameters, the profiles are grouped in latitude zones, and the
303 mean and standard values of the atmospheric parameters are computed at each altitude level. The latitude groups are
304 summarized in Table 2. They have been chosen such that each group is statistically meaningful, i.e. that the number of
305 measurements is statistically reliable. For that reason, the first two zones are 30 degrees large while above 60°, zones of
306 10° were considered. In the present study, symmetry in terms of latitude with respect to the Equator is assumed, as it has
307 been observed by the VIRTIS instrument on board Venus Express [Migliorini *et al.*, 2012]. Global circulation modeling
308 predicts different wind and dynamics characteristics at the AM and PM sides of the terminator, that should be reflected
309 on density and temperature [Brecht, 2011]. However, symmetry between the morning and evening sides of the
310 terminator has been assumed in this study. This is justified by the fact that no systematic variations were observed either
311 in terms of latitude or terminator side on the measurements investigated in this dataset. This observation will be verified
312 in the future, considering new measurements campaigns. These hypotheses also ensure a statistically valid number of
313 measurements in each latitude region, which would not be the case if AM and PM measurements had to be considered
314 separately.

315

316 Table 2

317

318 VAST profiles for the different latitude regions are shown in Figure 9 for the CO₂ density and in Figure 10 to 14 for the
319 temperature. In terms of CO₂ density, below 100 km VAST CO₂ density agrees with the Zasova model [Zasova *et al.*,
320 2007; Zasova *et al.*, 2006]. Above 100 km, VAST agrees better with the Keating model, based on the Pioneer Venus
321 measurements [Keating *et al.*, 1980], whereas the temperature profiles are more consistent with the VTS3 model [Hedin
322 *et al.*, 1983]. VAST CO₂ densities have also been compared to those from von Zahn *et al.* (1980) in Figure 9. The
323 Pioneer Venus data were taken in the Southern hemisphere (-40°), close to the morning side of the terminator. There is a
324 good agreement from 100 to 130 km between the two data sets, but the Pioneer Venus data do not show the same
325 pronounced gradient change at 125 km nor the lower densities observed by the SOIR instrument above. The difference
326 between the two models, VTS3 and Keating, at higher altitudes is striking and clearly indicates that observations
327 performed by SOIR can be accounted for. The region between 120 and 140 km in which SOIR observations show a
328 pronounced change in curvature is not reproduced by any of the models, but this is no surprise since both the Keating and
329 the VTS3 models have been considered to start at 140 km, even if they return data down to 100 km. Nothing on this
330 region can therefore be ascertained from these models. However, similar changes of curvature of the CO₂ density in the

331 120 to 130 km of altitude are reproduced by Venus Thermospheric General Circulation Model (VTGCM) [*Bougher et*
332 *al.*, 1988; *Brecht et al.*, 2012; *Parkinson*, 2012]. Detailed simulations by VTGCM of typical terminator profiles, such as
333 those observed by SOIR, are currently under investigation.

334

335 Figure 9

336

337 In the next paragraphs, VAST temperature profiles are compared to previously measured temperature profiles, for each
338 latitude region. These are represented by different symbols in each plot of Figure 10 to 14, where they are placed in the
339 corresponding figure when the latitude of the measurement is available. For most of the ground based measurements, the
340 exact latitude and local solar time are not always specified or well characterized, as the field of view of the observing
341 instrument often covers a wide range of latitudes and longitudes. For this reason, they are compared to mid-latitudes
342 measurements only – from the Equator to 60°, in Figure 10 to 12. For all latitude regions, there is a good agreement in
343 the 1 mbar to 10⁻² mbar region (95 to 102 km region). The SOIR temperature measurements have recently been
344 compared to simultaneous ground base measurements by *Sonnabend et al.* (2012), and a good agreement (10 K) between
345 the two data sets was reported.

346

347 Figure 10

348

349 For the latitudinal region extending from the Equator to 30° of latitude (see Figure 10), there is a good agreement with
350 the SPICAV measurements that were obtained close to the sub-solar point in the 10⁻² to 10⁻⁵ mbar region (102 to
351 125 km) [*Bertaux et al.*, 2007b]. However, their inversion observed at 90-100 km is not present in the SOIR
352 temperatures profiles, which show an inversion at a slightly higher altitude, around 110 km, with colder temperatures,
353 around 190 K. The comparison with the Pioneer Venus descent probe from [*Seiff et al.*, 1980] is also in very good
354 agreement in the 10 to 10⁻³ mbar region (80 to 105 km). Also, there is a good agreement with ground based
355 measurements [*Clancy et al.*, 2003; *Clancy et al.*, 2012; *Lellouch et al.*, 1994], but also with the Pioneer Venus OIR
356 measurement [*Taylor et al.*, 1980], with the measurements remaining within the SOIR error bars up to a pressure level of
357 10⁻² mbar (100 km). However, there is a disagreement at higher altitudes, above the 10⁻² mbar pressure level (100 km),
358 where SOIR derives warmer temperatures. There is a source of uncertainty linked to the local solar time of some of the
359 ground based measurements. Moreover they were not taken necessarily at the Venus terminator, like the SOIR
360 measurements presented in this study. For example, all the SPICAV measurements [*Bertaux et al.*, 2007b] were obtained
361 from the deep night side of Venus. This might explain some of the observed discrepancies. Also, there is a large
362 temperature variation of the SOIR data set (50 K) at all latitudes. Finally, a warm layer is observed at higher altitude

363 corresponding to the 10^{-7} mbar pressure level or 140-145 km, which cannot be corroborated nor dismissed by existing
364 ground-based or space observations.

365

366 Figure 11

367

368 In the 30° to 60° latitude region presented in Figure 11, the SOIR mean profile is in agreement with some of the
369 literature profiles in the 1 to 0.1 mbar region (80 to 90 km). At the 0.1 mbar level, the temperature inversion reported in
370 most of the literature temperature profile is warmer (170 K) than the SOIR temperature (150 K) [Clancy *et al.*, 2003;
371 Clancy *et al.*, 2012; Lellouch *et al.*, 1994; Seiff *et al.*, 1980; Taylor *et al.*, 1980]. Above the 0.1 mbar level, a warm layer
372 is observed in the SOIR data around 10^{-3} mbar (112 km) with a temperature of 225 K, decreasing to a cold layer reaching
373 100 K at 10^{-5} mbar (125 km). Few of the literature profiles reported here show such a structure at these latitudes, except
374 for the SPICAV measurements which also reach such cold temperature values – in the deep Venus nightside, but they do
375 not show the same gradient neither coincide with the VAST profile [Clancy *et al.*, 2012]. VAST also agrees with the
376 Keating and Zasova models at low altitude, at 1 mbar or 85 km. There is a warm layer observed at higher altitude
377 corresponding to the 10^{-7} mbar pressure level or 140-145 km.

378

379 Figure 12

380

381 In the 60° to 70° Lat region (Figure 12), the warm layer observed in VAST at the $5 \cdot 10^{-3}$ mbar pressure level (105 km) is
382 colder (150 K) than in the 30° to 60° Lat region (180 K). VAST agrees with the ground based measurements [Clancy *et al.*
383 *et al.*, 2012; Lellouch *et al.*, 1994; Taylor *et al.*, 1980] and VeRa on board Venus Express measurements [Piccialli *et al.*,
384 2012] from the 1 mbar to 10^{-2} mbar pressure level (85 km to 100 km). At higher altitude, VAST shows a warm layer that
385 is not reported in the literature at the $5 \cdot 10^{-3}$ mbar pressure level (105 km). Comparing the VAST temperature profile
386 with the measurements obtained by [Clancy *et al.*, 2011] on the dayside and nightside of the Venus, the VAST
387 minimum-maximum structure is also observed, such as in their nighttime profile, but VAST presents an intermediate
388 regime between these two, with the maximum temperature located at higher altitude. The cold layer at the 10^{-5} mbar
389 pressure level (125 km) reaches temperatures down to 90 K, which were already measured from ground based [Lellouch
390 *et al.*, 1994] but at slightly lower altitude (10^{-4} mbar or 115 km) and on the Venus nightside. A warm layer is also
391 observed at higher altitude, 10^{-7} mbar or 140-145 km.

392

393 Figure 13

394

395 In the 70° to 80° latitude region, see Figure 13, the profiles are not compared to ground based measurements as the
396 latitude is too high, and ground based measurements have rather difficult access to these regions. Only the Keating and
397 Zasova models are presented. There is a warm layer at the $8 \cdot 10^{-3}$ mbar (105 km) level, which located at slightly lower
398 altitude than in the lower latitude regions. The cold layer at the 10^{-5} mbar (125 km) level is observed at the same pressure
399 level than in the other latitude regions, such as the warm layer located at 10^{-7} mbar level (140-145 km).

400

401 Figure 14

402

403 In the 80° to 90° region (Figure 14), the SOIR temperature profile agrees almost perfectly, within the error bars, with the
404 VeRa profile in the same latitude region [Tellmann *et al.*, 2009; Piccialli *et al.*, 2012], between 80 mbar and $3 \cdot 10^{-2}$ mbar
405 (70 to 95 km). Above, the warm layer is located at 10^{-2} mbar (100 km), so slightly at lower altitude than at lower
406 latitude, and the cold layer remains at the same pressure level (10^{-5} mbar or 125 km) and with the same temperature
407 minimum (~ 90 K) [Clancy *et al.*, 2012]. The warm layer observed at higher altitude is also located at a pressure level of
408 10^{-7} mbar (140-145 km).

409

410 The temperature decrease at the 10^{-3} to 10^{-5} mbar level region (110-130 km region) above the temperature maximum at
411 $5 \cdot 10^{-3}$ mbar level (100-120 km region) has already been observed by the SPICAV instrument [Bertaux *et al.*, 2007b],
412 and is observed on the Venus night side by measurements made by VeRa [Piccioni *et al.*, 2009] on board Venus
413 Express. The striking temperature minimum observed in the SOIR dataset is always located at the same pressure level
414 (10^{-5} mbar or 125 km), and reaches the same range of temperature (80-90 K).

415

416 A comparison with night side measurements obtained at different latitudes by VIRTIS on board Venus Express
417 [Migliorini *et al.*, 2012] has also been carried out. The measurements described in that paper never occurred at the
418 terminator, but at close local solar times (5.00 AM and 5.00 PM). The temperatures are close to the SOIR data presented
419 here, between 190 and 220 K in the 1 to 100 mbar region. Also, there is a good agreement (less than 5 K difference) with
420 the temperature profiles obtained by the VeRa instrument on board VEX at pressure levels between 100 and 0.1 mbar
421 presented in Piccialli *et al.* (2012) at four different latitudes (-7° , -33° , -68° and -85° , see Figure 10 to 14) in the
422 Southern hemisphere, even if these measurements were obtained either at midnight or at 10.00 AM.

423 **8 Conclusions**

424 An atmospheric model for the Venus mesosphere and lower thermosphere has been constructed based on SOIR
425 observations. Those correspond to solar occultation probing the altitude range from 70 to 170 km with a high vertical
426 resolution. They cover a wide range of latitudes, but correspond all to the 6.00 AM and 6.00 PM local times, as all

427 observations are performed at the terminator, either on its morning or its evening side. A subset of 59 observations has
428 been considered, from which information on the CO₂ density and the temperature were derived. VAST (Venus
429 Atmosphere from SOIR measurements at the Terminator) is defined on 5 latitudinal zones and provides CO₂ and total
430 densities and temperature as a function of the altitude. VAST has been compared to data available in the literature, and it
431 was shown that, in particular, the temperature profiles are in good agreement with literature data at the lowest altitudes.
432 A striking and never observed before cold layer is always observed at an altitude of 125 km.

433

434 VAST will be further refined and improved with other CO₂ measurements, already recorded or obtained during future
435 dedicated SOIR observation campaigns. The first step will be to analyze all observations already acquired allowing the
436 determination of the CO₂ density and temperature, but which were not considered in this study because they did not
437 cover a high enough altitude range. The model will also be further expanded with the inclusion of the vertical profiles of
438 various trace gases observed by SOIR, such as H₂O and HDO [*Fedorova et al.*, 2008], HCl, HF, CO [*Vandaele et al.*,
439 2008], SO₂ [*Belyaev et al.*, 2008; *Belyaev et al.*, 2012].

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445 Agency (ESA, PRODEX program, contracts C 90268, 90113, and 17645).

446

447 **10 Tables captions**

448 Table 1: List of the 59 occultations studied in this work. OS is the occultation season number, LST is the mean local
449 solar time of the measurement in hour, Lat is the mean latitude in degrees and Lon is the mean longitude in degrees. The
450 decimal after the orbit number represents the measurement number of the day. The orders used to derive information on
451 the CO₂ density and temperature profiles are also indicated. The orders 121 (2704 to 2727 cm⁻¹), 134 (2995 to
452 3020 cm⁻¹), 171 (3822 to 3854 cm⁻¹), 176 to 180 (3933 to 4057 cm⁻¹) give the CO₂ density at low altitudes (70 to 90 km).
453 Mid-altitude (80 to 120 km) CO₂ density is derived from orders 112 (2503 to 2525 cm⁻¹), 117 and 118 (2615 to 2637 cm⁻¹),
454 141 (3151 to 3178 cm⁻¹), 147 to 150 (3285 to 3381 cm⁻¹). The orders 155 and 156 (3464 to 3516 cm⁻¹) give
455 information on the CO₂ density at high altitudes (100 to 140 km), and orders 160 to 168 (3576 to 3787 cm⁻¹) at very high
456 altitudes (120 to 170 km).

457

458 Table 2 : Characteristics of the different latitude zones defined in VAST. Each group is characterized by a minimum and
459 maximum latitude. The numbers of orbits, and the corresponding minimum and maximum probed altitudes of each
460 group, are also indicated.

461

462 **11 Figure captions**

463 Figure 1 : Working principle of the iterative inversion scheme. At each iteration, as a first step, the eight sets of spectra
464 are inverted using the Asimat algorithm without fitting the temperature. Then, the eight independent CO₂ density
465 profiles are recombined into one single smooth profile. The last step consists in calculating the temperature profile using
466 the hydrostatic law and the CO₂ VMR from the Keating and Zasova models [Keating *et al.*, 1985; Zasova *et al.*, 2007;
467 Zasova *et al.*, 2006]. The convergence criteria are used to check if more iterations are needed.

468

469 Figure 2: Illustration of the convergence of the atmospheric parameters iterative procedure for orbit 687.1. The density
470 convergence (left panel) and the temperature convergence (center panel) are presented. The parameters that are plotted
471 for both variables are the integrated difference between two successive profiles. The error bars represent the average
472 density and temperature error. The residual mean square summed on all the fitted spectra that are considered to build the
473 profiles at each step of the occultation is shown in the right panel. The convergence criterion is set to 0.3 for the density
474 logarithm and 7 K for the temperature, and is shown as the thick line in the left and center panels. The convergence
475 occurs after three steps. The further steps are displayed to show the stabilization, within the error bar. After 2 steps, the
476 RMS is getting lower, see right panel.

477

478 Figure 3: Convergence of the CO₂ density (left panel) and temperature profiles (right panel) for orbit 687.1 during the
479 iterative procedure. The Zasova and the Keating models are presented as the black lines [Keating *et al.*, 1985; Zasova
480 *et al.*, 2007; Zasova *et al.*, 2006]. The convergence occurred after 3 steps, both in terms of temperature and density (see
481 Figure 2). The converged profiles are the thick green profiles. One step further is also displayed to show the stabilization
482 of the derived profiles.

483

484 Figure 4: Influence of the initial density value on the derived temperature profile presented for orbit 687.1. The
485 temperature profiles are calculated using the hydrostatic law for different initial density values, ranging from
486 $\rho_{CO_2}(z_0) - 2 \cdot \sigma$ and $\rho_{CO_2}(z_0) + 2 \cdot \sigma$, with σ the density error on the maximum altitude point, around the density
487 derived at the maximum altitude, $\rho_{CO_2}(z_0)$. The gray zone is the calculated error on the profile at convergence. It has
488 typical values of 1 to 20 K. All the profiles starting from different initial temperatures align below two scale heights,
489 which corresponds to 10 km at this altitude, and remain within the error bar of the converged profile at lower altitudes.

490

491 Figure 5: Example of two sets of profiles. The first one is a North polar measurement (orbit 687.1, left panels) and the
492 second one an equatorial measurement (orbit 1567.1, right panels). The main difference between the two profiles resides

493 in the vertical resolution and the number of measurement points, which are lower when going to Southern latitudes. In
494 the CO₂ density Panels (top panels), the points are the retrieved individual density values and the black line is the
495 averaged profile. The error ranges between 1 and 5 %. In the Temperature Panels (bottom panels), the gray envelope is
496 the error on the temperature. Typical temperature error values are 1 to 20 K.

497

498 Figure 6: Localization of the orbit dataset considered in the study. The Top Panel gives the local solar time of the
499 measurement as a function of the latitude. The Bottom Panel gives the orbit number as a function of the latitude of the
500 measurement. In this Panel, the occultation seasons considered are in gray, and their number is also indicated.

501

502 Figure 7: CO₂ density profiles (Left Panel) and CO₂ temperature profiles (Right Panel) of the orbit dataset considered in
503 this study. The inset panel gives the measurement latitude and the orbit number. The density profiles are given as a
504 function of the altitude, and the temperature profiles as a function of the total pressure, with the altitude given on the
505 right side as an indication. The color is the absolute latitude. High latitude measurements are reddish, while equatorial
506 are bluish. The black lines are the density and temperature values of the Keating [Keating *et al.*, 1980] (plain line) and
507 Zasova models [Zasova *et al.*, 2007; Zasova *et al.*, 2006], for different latitudes (dashed is 0°, dashed-dotted is 45° and
508 dotted is 90°).

509

510 Figure 8: Homopause altitudes derived from the SOIR measurements, together with the error values. The altitude of the
511 homopause derived from the SOIR data varies between 120 and 132 km.

512

513 Figure 9: VAST CO₂ density profiles in function of altitude for the different latitude regions defined in Table 2. The
514 thick line represents the VAST data; the gray envelope is its confidence range (ranging between 1 and 5 %). The Keating
515 model [Keating *et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black line [Zasova
516 *et al.*, 2007; Zasova *et al.*, 2006] and the VTS3 model [Hedin *et al.*, 1983] by the thin dashed line. The density profile
517 obtained Pioneer Venus is also given [von Zahn *et al.*, 1980].

518

519 Figure 10: VAST temperature profiles in function of the total pressure in the 0 to 30° latitude region. The altitude levels
520 are given as indications in the boxes. The thick line represents the VAST data; the gray envelope is its confidence range.
521 The Keating model [Keating *et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black
522 line [Zasova *et al.*, 2007; Zasova *et al.*, 2006] and the VTS3 model [Hedin *et al.*, 1983] by the thin dashed line. Different
523 measurements from the literature, from ground-based or spacecraft measurements, are also indicated [Bertaux *et al.*,

524 2007b; *Clancy et al.*, 2003; *Clancy et al.*, 2012; *Lellouch et al.*, 1994; *Piccialli et al.*, 2012; *Seiff et al.*, 1980; *Taylor et*
525 *al.*, 1980 and references therein]. The local solar times are given when they are known.

526

527 Figure 11: VAST temperature profiles in function of total pressure in the 30 to 60° latitude region. The altitude levels are
528 given in the boxes. The thick line represents the VAST data; the gray envelope is its confidence range. The Keating
529 model [*Keating et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black line [*Zasova*
530 *et al.*, 2007; *Zasova et al.*, 2006] and the VTS3 model [*Hedin et al.*, 1983] by the thin dashed line. Different
531 measurements from the literature, from ground-based or spacecraft measurements, are also indicated [*Bertaux et al.*,
532 2007b; *Clancy et al.*, 2003; *Clancy et al.*, 2012; *Lellouch et al.*, 1994; *Piccialli et al.*, 2012; *Seiff et al.*, 1980; *Taylor et*
533 *al.*, 1980 and references therein]. The local solar times are given when they are known.

534

535 Figure 12: VAST temperature profiles in function of total pressure in the 60 to 70° latitude region. The altitude levels are
536 given in the boxes. The thick line represents the VAST data; the gray envelope is its confidence range. The Keating
537 model [*Keating et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black line [*Zasova*
538 *et al.*, 2007; *Zasova et al.*, 2006] and the VTS3 model [*Hedin et al.*, 1983] by the thin dashed line. Different
539 measurements from the literature, from ground-based or spacecraft measurements, are also indicated [*Clancy et al.*,
540 2003; *Clancy et al.*, 2012; *Lellouch et al.*, 1994; *Piccialli et al.*, 2012; *Taylor et al.*, 1980]. The local solar times are
541 given when they are known.

542

543 Figure 13: VAST for the temperature profiles in function of total pressure in the 70 to 80° latitude region. The altitude
544 levels are given in the boxes. The thick line represents the VAST data; the gray envelope is its confidence range. The
545 Keating model [*Keating et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black line
546 [*Zasova et al.*, 2007; *Zasova et al.*, 2006] and the VTS3 model [*Hedin et al.*, 1983] by the thin dashed line.

547

548 Figure 14: VAST temperature profiles in function of total pressure in the 80 to 90° latitude region. The altitude levels are
549 given in the boxes. The thick line represents the VAST data; the gray envelope is its confidence range. The Keating
550 model [*Keating et al.*, 1980] is illustrated by the thin dash-dotted line, the Zasova model by the plain black line [*Zasova*
551 *et al.*, 2007; *Zasova et al.*, 2006] and the VTS3 model [*Hedin et al.*, 1983] by the thin dashed line. Different
552 measurements from the literature, from ground-based or from spacecraft, are also indicated [*Piccialli et al.*, 2012;
553 *Tellmann et al.*, 2009]. The local solar times are given when they are known.

554

555

Orbit	Date	OS	LST	Lat	Lon	CO ₂ orders			
						-	-	-	-
-	DD/MM/YY	-	Hour	Deg	Deg	-	-	-	-
35.1	26/05/06	1	5.205	76.7	306	166	167		
361.1	17/04/07	4	6.862	79.8	207	112	147	149	166
593.1	05/12/07	6	6	2	215	112	147	149	162
597.2	09/12/07	6	6	64	222	111	149	161	119
667.1	17/02/08	7	18	80	258	112	149	166	
669.1	19/02/08	7	19	81	262	112	149	166	
671.1	21/02/08	7	19	82	266	112	149	166	
674.1	24/02/08	7	19	84	271	112	149	166	
675.1	25/02/08	7	19	84	272	112	149	166	
677.1	27/02/08	7	19	85	274	112	149	166	
679.1	29/02/08	7	20	86	273	112	149	166	
681.1	02/03/08	7	20	87	270	112	149	166	
684.1	05/03/08	7	22	88	254	112	149	166	
685.1	06/03/08	7	23	88	245	112	149	166	
686.1	07/03/08	7	24	88	235	112	149	166	
687.1	08/03/08	7	1	88	225	112	149	166	
688.1	09/03/08	7	1	88	216	112	149	166	
689.1	10/03/08	7	2	87	209	149	166		
690.1	11/03/08	7	3	87	204	149	166		
691.1	12/03/08	7	3	86	201	149	166		
693.1	14/03/08	7	4	86	197	134	149	166	
695.1	16/03/08	7	4	84	197	149	166		
697.1	18/03/08	7	4	83	198	149	166		
699.1	20/03/08	7	5	82	201	149	166	178	
703.1	24/03/08	7	5	78	208	112	149	166	
705.1	26/03/08	7	5	76	212	149	166		
706.1	27/03/08	7	5	75	214	149	166		
709.1	30/03/08	7	5	70	220	112	149	166	
710.1	31/03/08	7	6	68	222	117	118	149	166
710.2	31/03/08	7	6	18	217	149	166	121	171
711.1	01/04/08	7	6	60	220	121	149	166	171
711.2	01/04/08	7	6	26	220	149	166	121	171
1121.1	16/05/09	11	19	83	212	161	162		
1124.1	19/08/09	11	18	-44	234	121	149	166	171
1126.1	21/08/09	11	18	-52	240	149	166		
1128.1	23/08/09	11	18	-61	248	149	161	162	166
1130.1	25/08/09	11	18	-69	255	149	166		
1132.1	27/08/09	11	18	-77	264	149	166		
1137.1	01/06/09	11	1	88	166	160	163	164	165
1139.1	03/06/09	11	7	-75	91	121	149	166	171
1141.1	05/06/09	11	6	-67	100	149	166		
1154.1	18/06/09	11	6	-8	144	149	161	162	166
1252.1	24/09/09	12	5	-74	95	149	164	165	166
1254.1	26/09/09	12	6	-65	98	149	166	167	168
1256.1	28/09/09	12	6	-57	103	121	149	166	171
1381.2	31/01/10	13	6	13	126	149	160	165	
1461.1	21/04/10	14	17	85	205	117	148	160	165
1462.1	22/04/10	14	18	-48	192	148	156	166	
1464.1	24/04/10	14	16	86	224	111	150	160	165
1465.1	25/04/10	14	18	-60	200	111	149	160	166
1467.1	27/04/10	14	15	87	248	149	150	156	166
1469.1	29/04/10	14	19	-76	207	112	148	156	166
1480.1	10/05/10	14	9	87	31	141	150	156	165

1565.1	03/08/10	15	18	75	149	148	165	155	179
1567.1	05/08/10	15	18	-11	159	141	149	156	166
1581.1	19/08/10	15	21	87	155	148	165	155	179
1798.1	24/03/11	17	19	82	140	118	141	156	167
1802.1	28/03/11	17	19	85	142	141	156	167	
1804.1	30/03/11	17	20	86	142	118	141	156	167

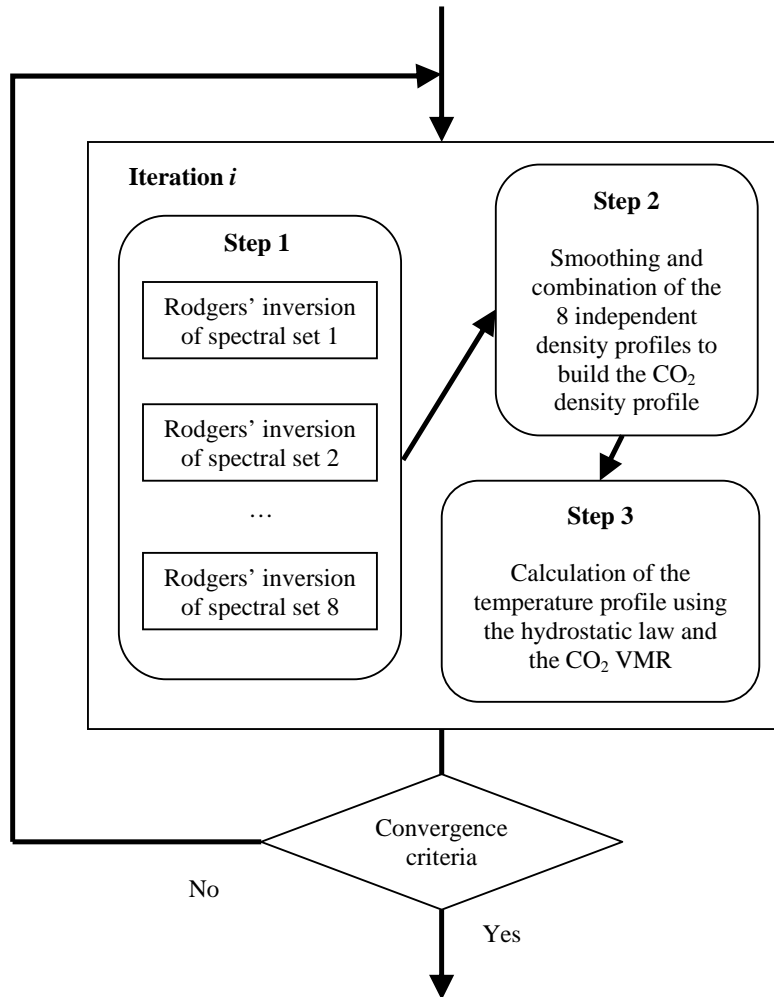
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Table 1

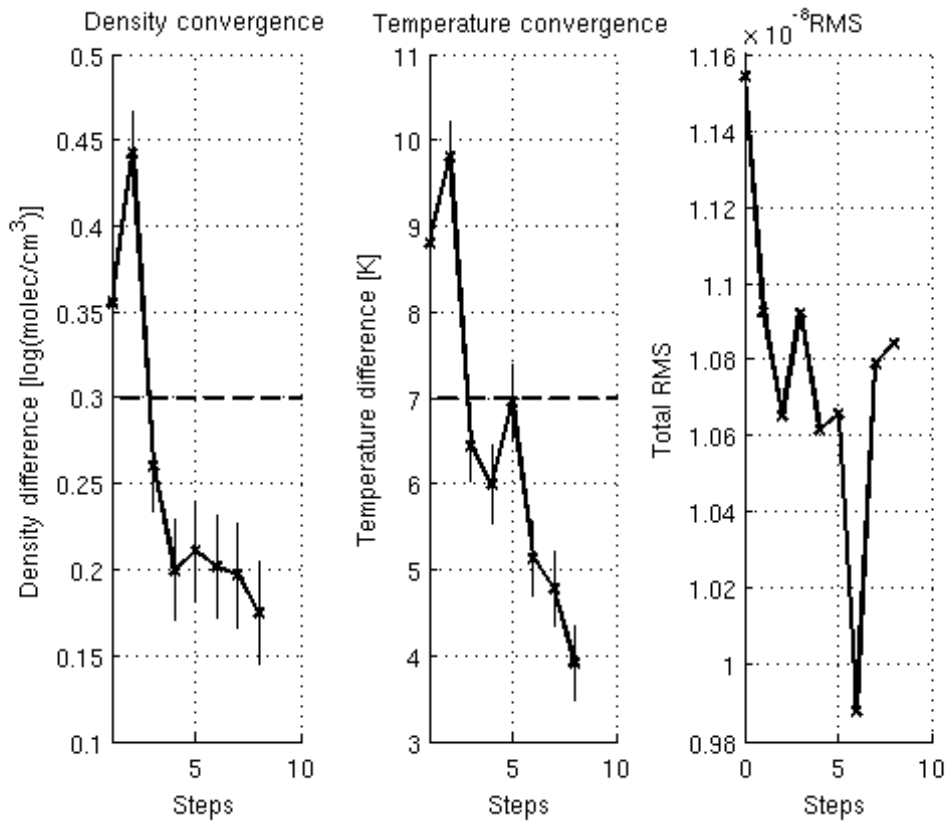
561

Latitude Min [deg]	Latitude Max [deg]	Number of orbits	Altitude Min [km]	Altitude Max [km]
0	30	6	82	165
30	60	4	78	161
60	70	9	81	168
70	80	11	73	168
80	90	29	62	161

562 Table 2
563

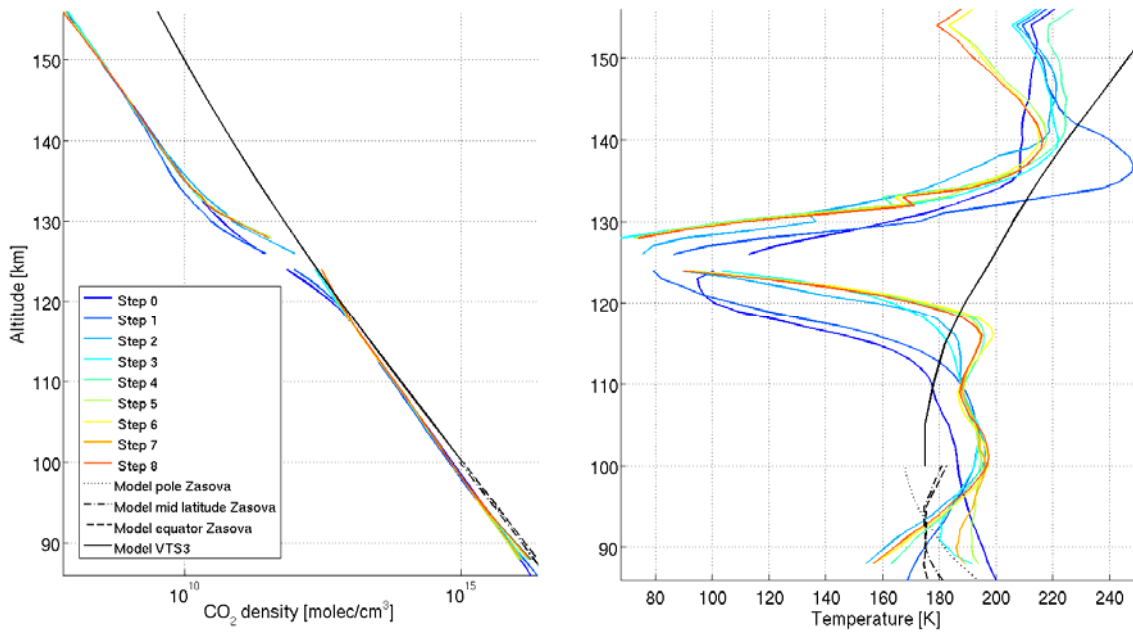


565
566 **Figure 1**
567
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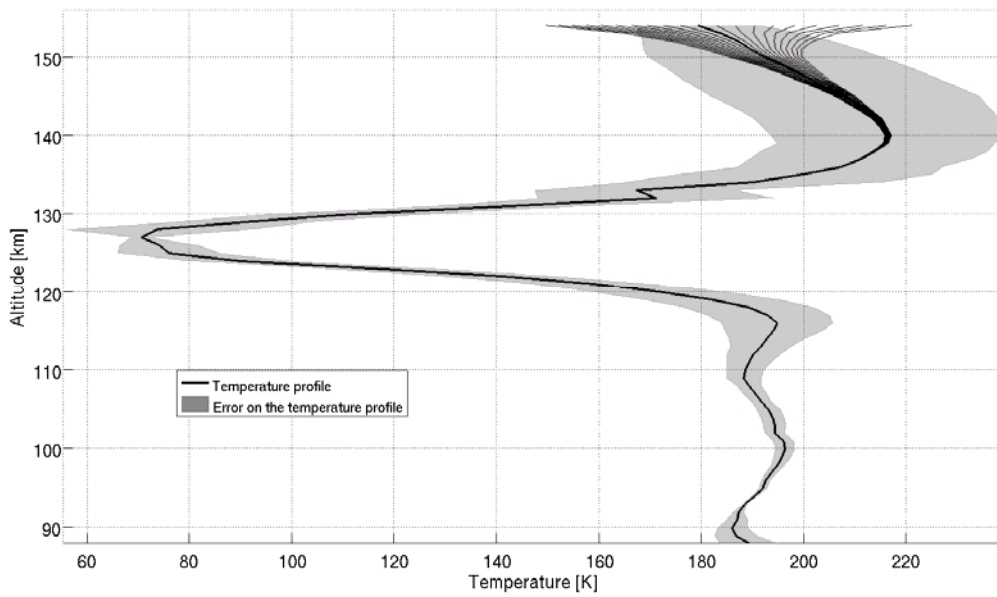
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Figure 2



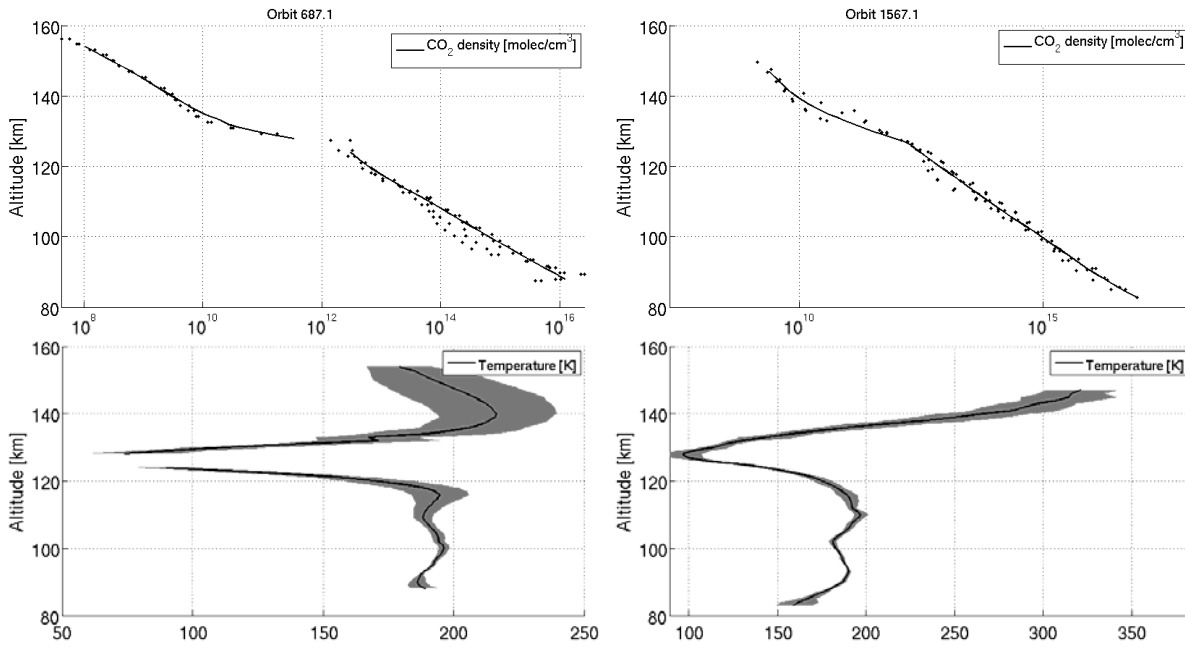
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Figure 3



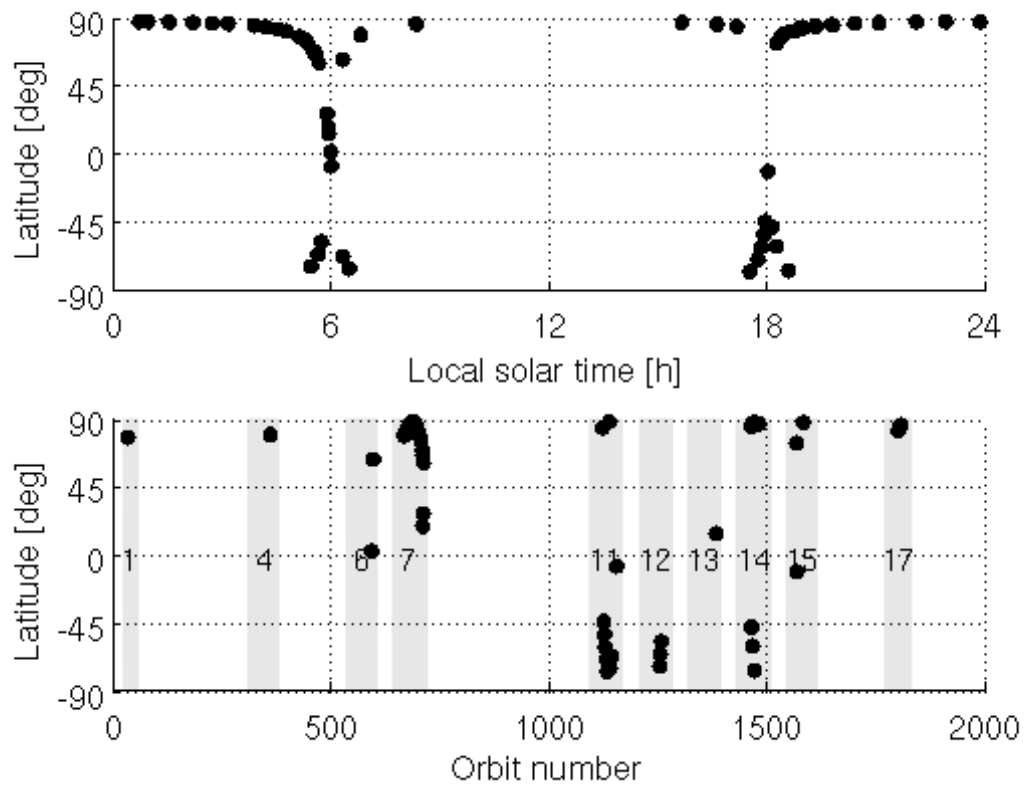
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Figure 4



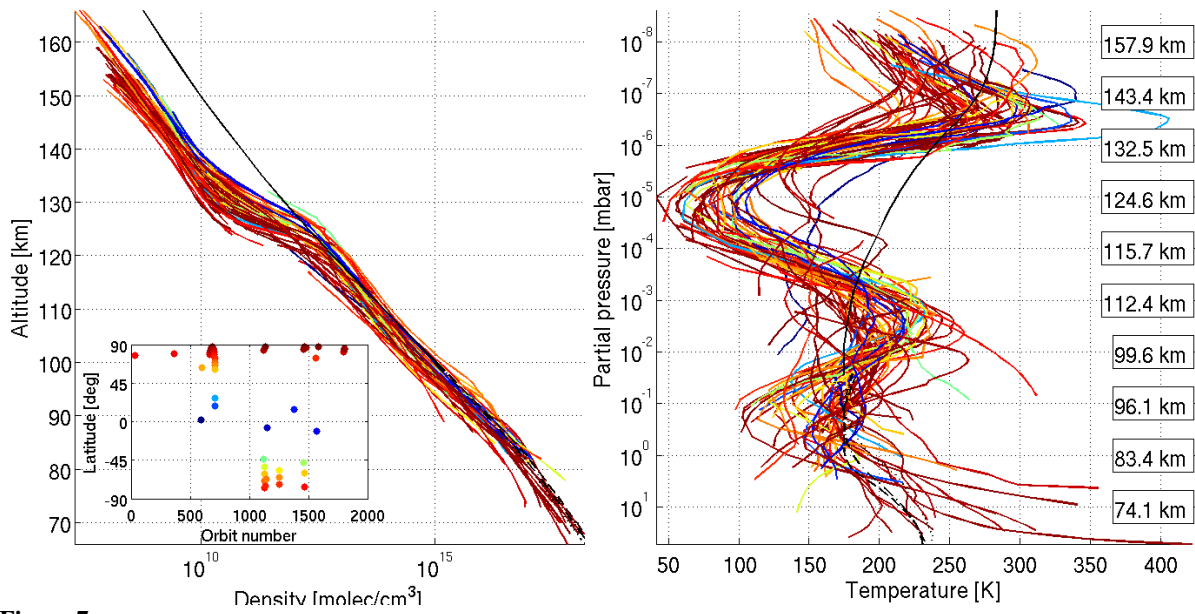
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Figure 5



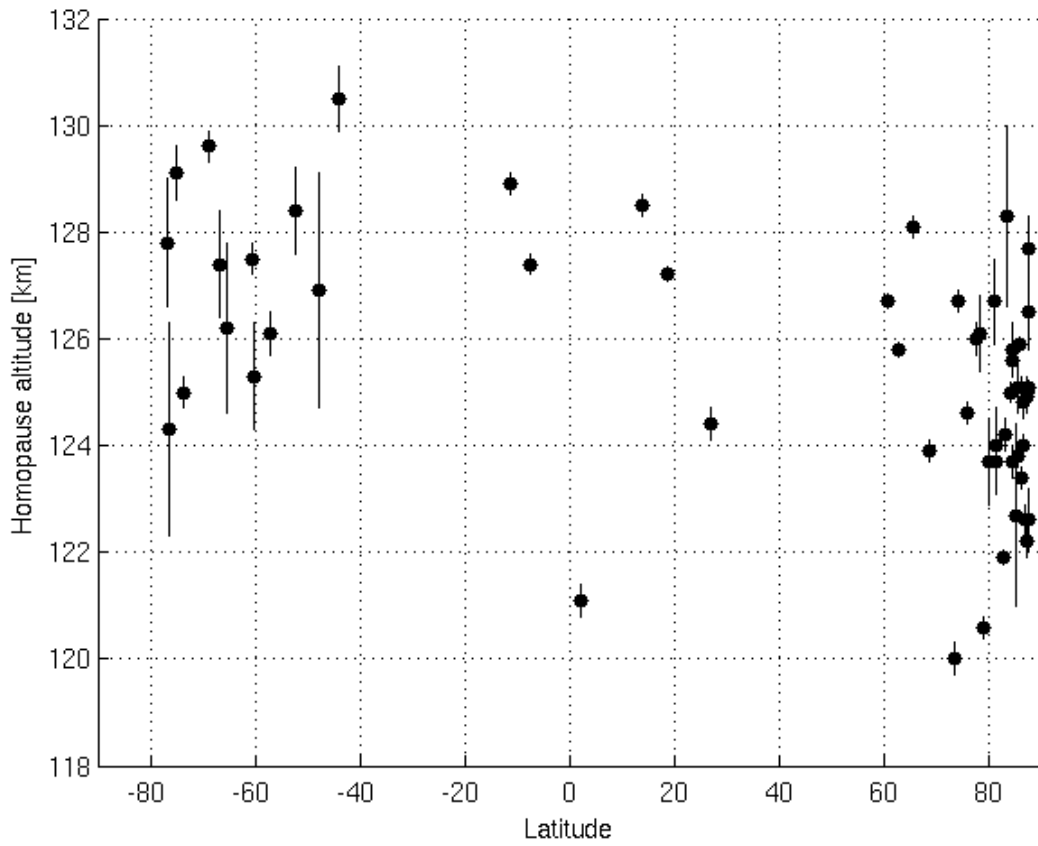
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Figure 6



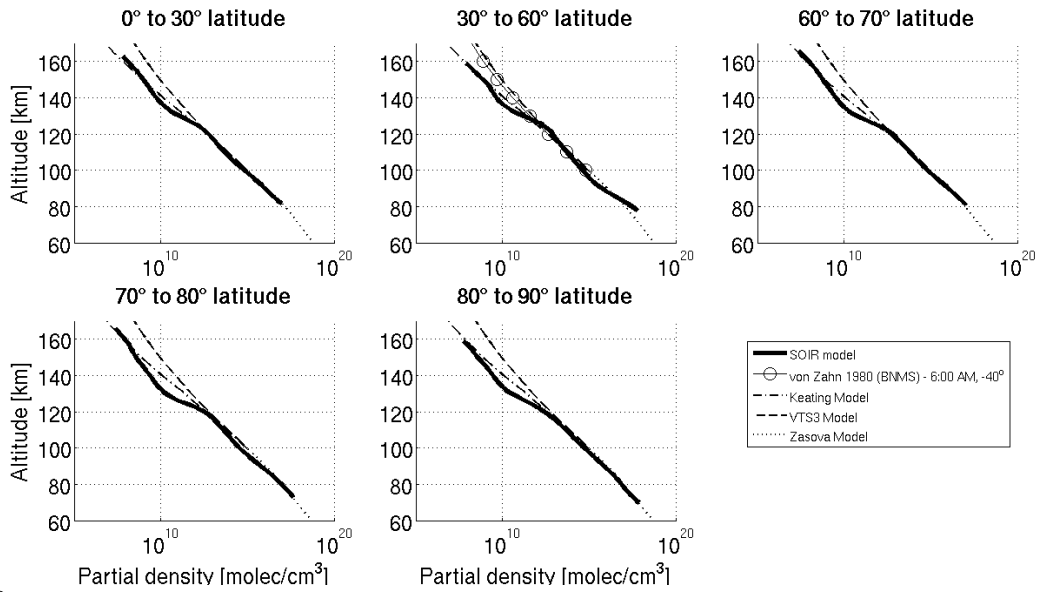
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Figure 7



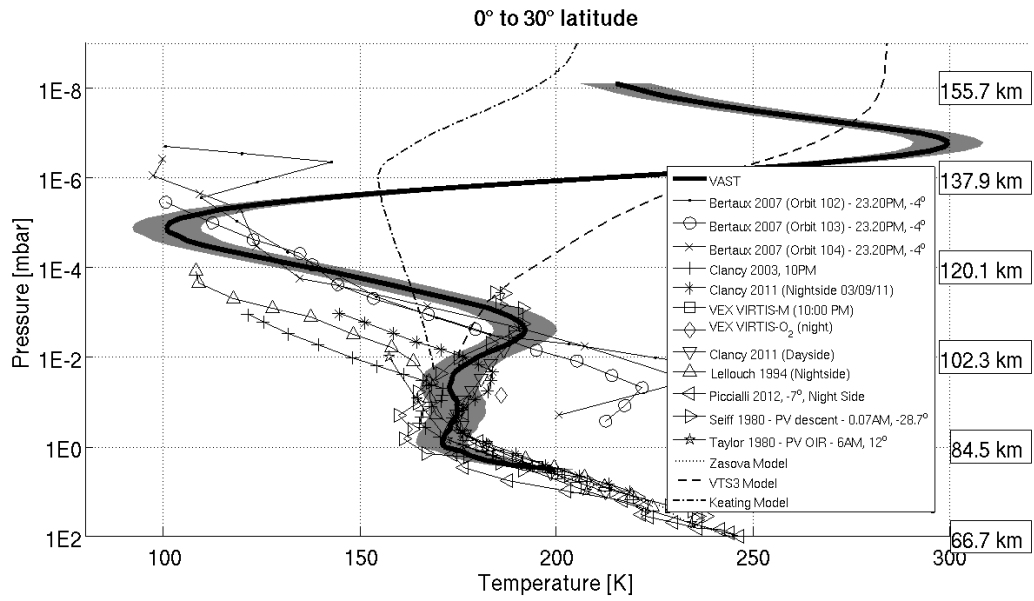
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Figure 8

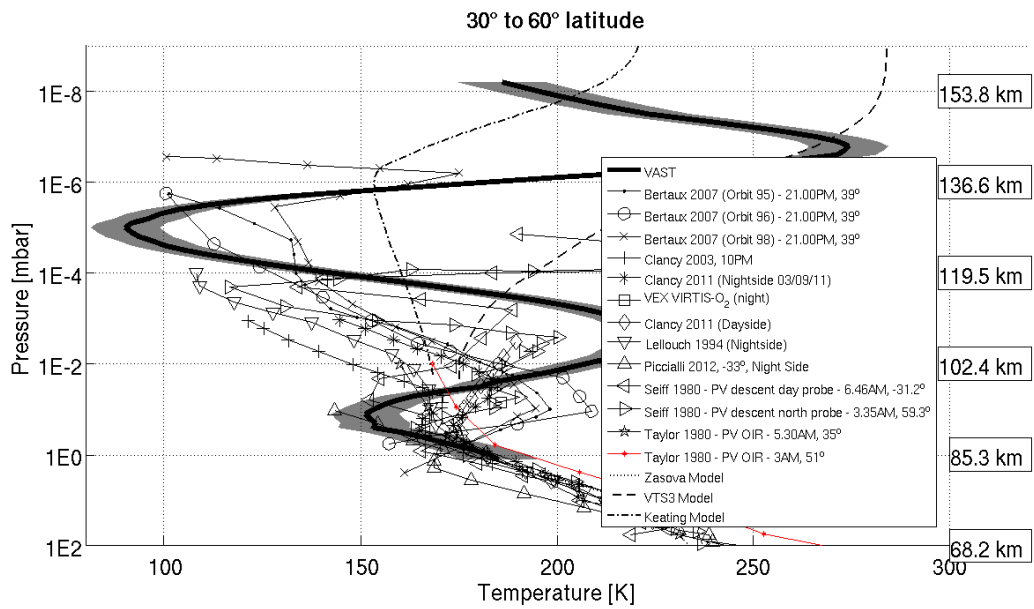


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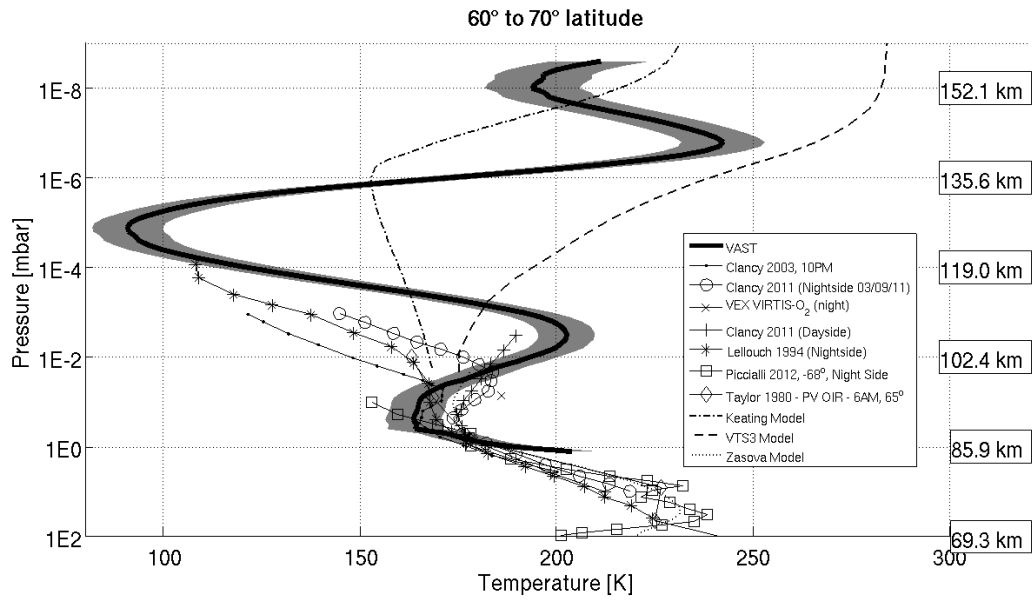
Figure 9



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594 **Figure 10**
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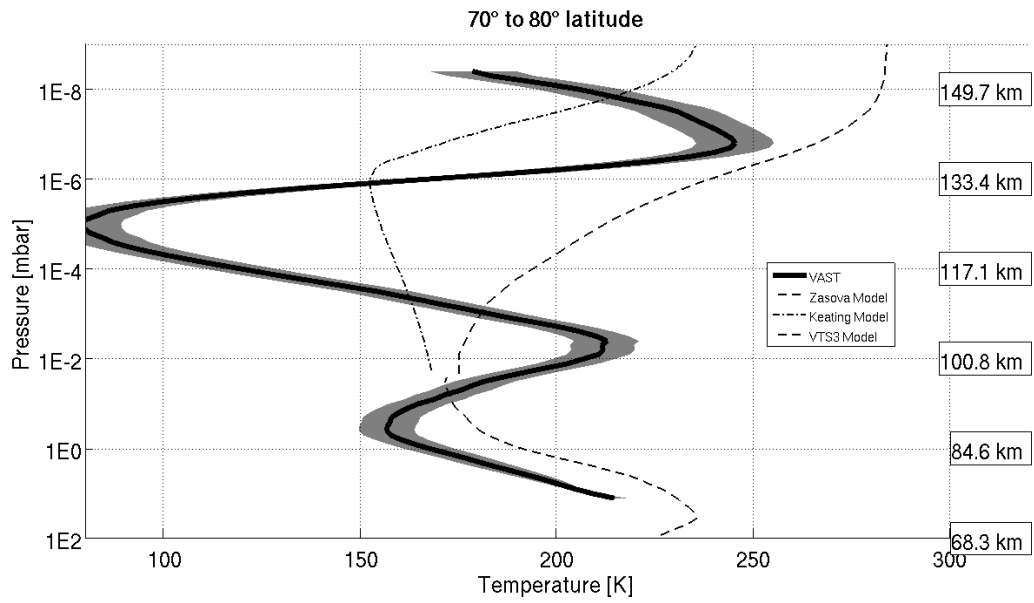


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597 **Figure 11**
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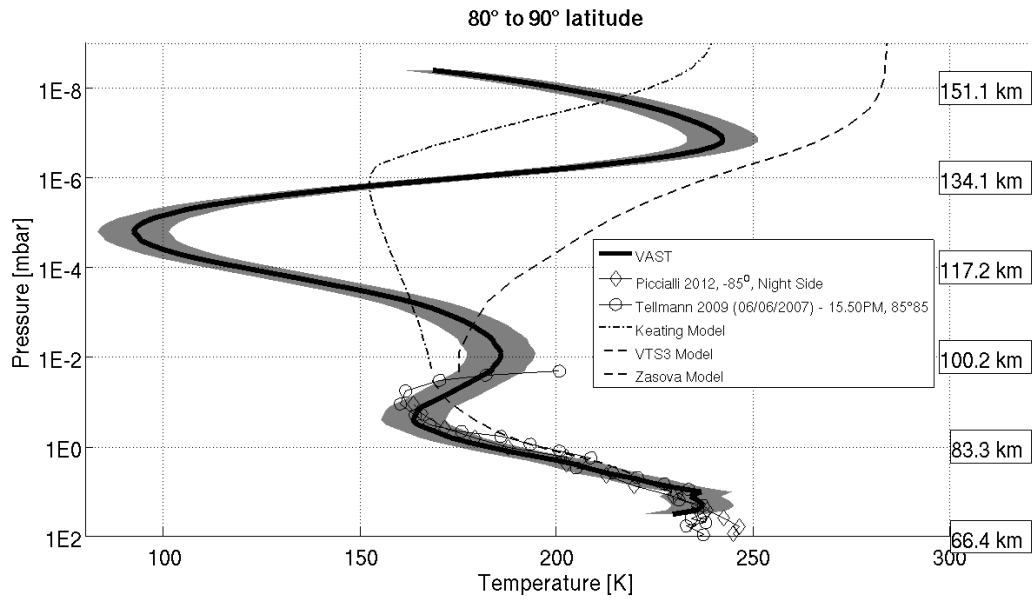
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Figure 12



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Figure 13



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606 **Figure 14**
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609 **14 References**

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